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## The silicon carbonyls revisited: on the existence of a planar $\text{Si}(\text{CO})_4$

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**Abstract** Recently, some works have focused attention on the reactivity of silicon atom with closed-shell molecules. Silicon may form a few relatively stable compounds with CO, i.e.  $\text{Si}(\text{CO})$ ,  $\text{Si}(\text{CO})_2$ ,  $\text{Si}[\text{C}_2\text{O}_2]$ , while the existence of polycarbonyl ( $n > 2$ ) silicon complexes has been rejected by current literature. In this paper, the reaction of silicon with carbonyl has been reinvestigated by density functional calculations. It has been found that the tetracoordinated planar  $\text{Si}(\text{CO})_4$  complex is thermodynamically stable. In  $\text{Si}(\text{CO})$ , silicon carbonyl, and  $\text{Si}(\text{CO})_2$ , silicon dicarbonyl, the CO are datively bonded to Si;  $\text{Si}(\text{CO})_4$ , silicon tetracarbonyl, may be viewed as a resonance between the extreme configurations  $(\text{CO})_2\text{Si} + 2\text{CO}$  and  $2\text{CO} + \text{Si}(\text{CO})_2$ ; while  $\text{Si}[\text{C}_2\text{O}_2]$ , *c*-silicodiketone, is similar to the compounds formed by silicon and ethylene. A detailed orbital analysis has shown that the Si bonding with two CO is consistent with the use of  $sp^2$ -hybridized orbitals on silicon, while the Si bonding with four CO is consistent with the use of  $sp^2d$ -hybridized orbitals on silicon, giving rise to a planar structure about Si.

**Keywords** Silicon carbonyl complexes · Planar tetracoordinated silicon · Density functional calculations

### 1 Introduction

In view of its numerous applications in the chemical (silicones) and electronic (single crystalline silicon) industries, silicon has been the subject of extended experimental and

theoretical investigations. Despite of that, many remarkable specific chemical behaviors (like the exceptional strength of its bond to fluorine and the formation of planar silyl amines) have remained partially unexplained and are still a matter for discussion.

In recent years, motivated by the fact that the thermal oxidation of single crystalline silicon in  $\text{O}_2$  ambient injects atomic silicon into the growing  $\text{SiO}_2$  [1,2], we have been attracted by the chemical behavior of atomic-like silicon in siloxanic networks. The major conclusion of our investigation was that silicon may behave as a weak bifunctional Lewis acid forming  $((-\text{O})_3\text{Si})_2\text{O} \rightarrow \text{Si} \leftarrow \text{O}(\text{Si}(\text{O}-)_3)_2$  adducts [3–5].

Atomic silicon may also be produced in such a state with relatively simple methods like sputtering or evaporation. These methods allow in principle an experimental study of the reaction of silicon with closed-shell molecules (for instance, as condensed film physisorbed at the surface of a substrate where the silicon atoms are being deposited) which are expected to behave to some extent as the siloxanic species mentioned above. In particular, the base nature of  $\text{NH}_3$ ,  $\text{H}_2\text{O}$  and CO (ordered for decreasing basic strength) would suggest for them the formation of  $\text{R} \rightarrow \text{Si} \leftarrow \text{R}$  adducts ( $\text{R} = \text{NH}_3, \text{H}_2\text{O}, \text{CO}$ ).

Recently, some works have focused attention on the reactivity of Si atom with unsaturated molecules (i.e.  $\text{C}_2\text{H}_2$  and  $\text{C}_2\text{H}_4$ ) [6,7].

This background has suggested to us a theoretical study of the expected species which may be formed in such conditions. Such analysis would be certainly useful while assessing the experiments and eventually rationalizing the results. We anticipate that while for  $\text{NH}_3$  and  $\text{H}_2\text{O}$  the calculations confirmed the expectations (P. Belanzoni, G. Giorgi, G.F. Cerofolini and A. Sgamellotti, unpublished), calculations provided surprising results for CO. Several papers have appeared in literature on  $\text{Si}(\text{CO})_n$  ( $n = 1, 2$ ) complexes, where the structures, properties and the bonding character for such silicon carbonyls have been investigated using different theoretical methods [8–10]. The existence of larger complexes was discarded on the basis of the weakness of the bond of

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the third CO to the complex and the absence of any reported experimental evidence for these species [8], while Si(CO) and Si(CO)<sub>2</sub> have been experimentally observed [11,12].

In this work, we, however, predict the thermodynamic stability of a planar tetracoordinated Si(CO)<sub>4</sub> complex.

## 2 Computational and methodological details

The calculations reported in this paper have been performed employing the Amsterdam Density Functional (ADF) program package [13–16]. The ADF code is characterized by the use of a density fitting procedure to obtain accurate Coulomb and exchange potentials in each self-consistent-field cycle in the solution of the one-electron Kohn–Sham equations, by accurate and efficient numerical integration [17] of the effective one-electron Hamiltonian matrix elements and by the possibility of freezing core orbitals. The molecular orbitals were expanded in a basis set of Slater-type orbitals (STOs) and the frozen core approximation was used for the evaluation of valence orbitals. The parameterization of electron gas data [18] by Vosko et al. [19] was employed in the local density approximation. Full geometry optimizations were performed within both spin unrestricted (open shell) and spin restricted (closed shell) approach including Becke’s gradient corrections [20] to the exchange part of the potential and Perdew’s gradient correction [21,22] to the correlation. For our systems, we considered different symmetries, i.e., C<sub>∞v</sub>, C<sub>2v</sub>, D<sub>∞h</sub>, D<sub>4h</sub>, D<sub>2h</sub>, T<sub>d</sub> and D<sub>2</sub>, in an attempt to find the “most” stable geometry in each series. For the model molecules, the Si, C, H, and O molecular orbitals were expanded in a triple- $\zeta$  STO basis set, adding as polarization functions one 3*d* STO for C and O, and one 3*d* plus one 4*f* STO for Si. The core orbitals (from 1*s* to 2*p* for Si, and 1*s* for O and C) were kept frozen. In order to analyze the silicon–carbonyl interaction energies, we used a method that is an extension of the well-known decomposition scheme of Morokuma [23,24]. The bonding energy is decomposed into a number of terms. The first term,  $\Delta E^0$ , is obtained from the energy of the wavefunction  $\Psi^0$  which is constructed as the antisymmetrized and renormalized product of the wavefunctions  $\Psi^A$  and  $\Psi^B$  of the fragments A and B from which the molecule is built up.  $\Delta E^0$ , which is called *steric repulsion*, consists of two components. The first is the *electrostatic interaction*,  $\Delta E_{\text{elstat}}$ , of the nuclear charges and unmodified electronic charge density of one fragment with those of the other fragment, both fragments being at their final positions. The second component is the so-called *exchange repulsion* or *Pauli repulsion*,  $\Delta E_{\text{Pauli}}$ , which is essentially due to the antisymmetry requirement of the total wavefunction. In addition to the steric repulsion term  $\Delta E^0$  there are attractive orbital interactions which enter when the wavefunction  $\Psi^0$  is allowed to relax to the fully converged ground state wavefunction of the total molecule,  $\Psi^{\text{AB}}$ . The energy lowering due to mixing of virtual orbitals of the fragments into the occupied orbitals is called *orbital interaction energy*,  $\Delta E_{\text{o.i.}}$ , that includes both the charge transfer and polarization interactions. This term, according

to the decomposition scheme proposed by Ziegler [25], may be broken up into contributions from the orbital interactions within the various irreducible representations of the overall symmetry group of the system. There is a third contribution to the total bonding energy ( $\Delta E = \Delta E^0 + \Delta E_{\text{o.i.}}$ ) in the frequent cases where the ground state wavefunctions  $\Psi^A$  and  $\Psi^B$ , at the equilibrium geometries of the free fragments, cannot be used to calculate  $\Delta E^0$ . The geometry of the free fragment is often different from the geometry of the fragments as it occurs in the overall molecule. Also, the ground electronic configuration of the free fragment may not be suitable for interaction with the other fragment. The energy required to prepare the fragments for interaction by changing the geometry and the electronic configuration is called *preparation energy*,  $\Delta E_{\text{prep}}$ . Thus the total bonding energy will be:

$$\Delta E = \Delta E_{\text{prep}} + \Delta E^0 + \Delta E_{\text{o.i.}} \quad (1)$$

## 3 Results

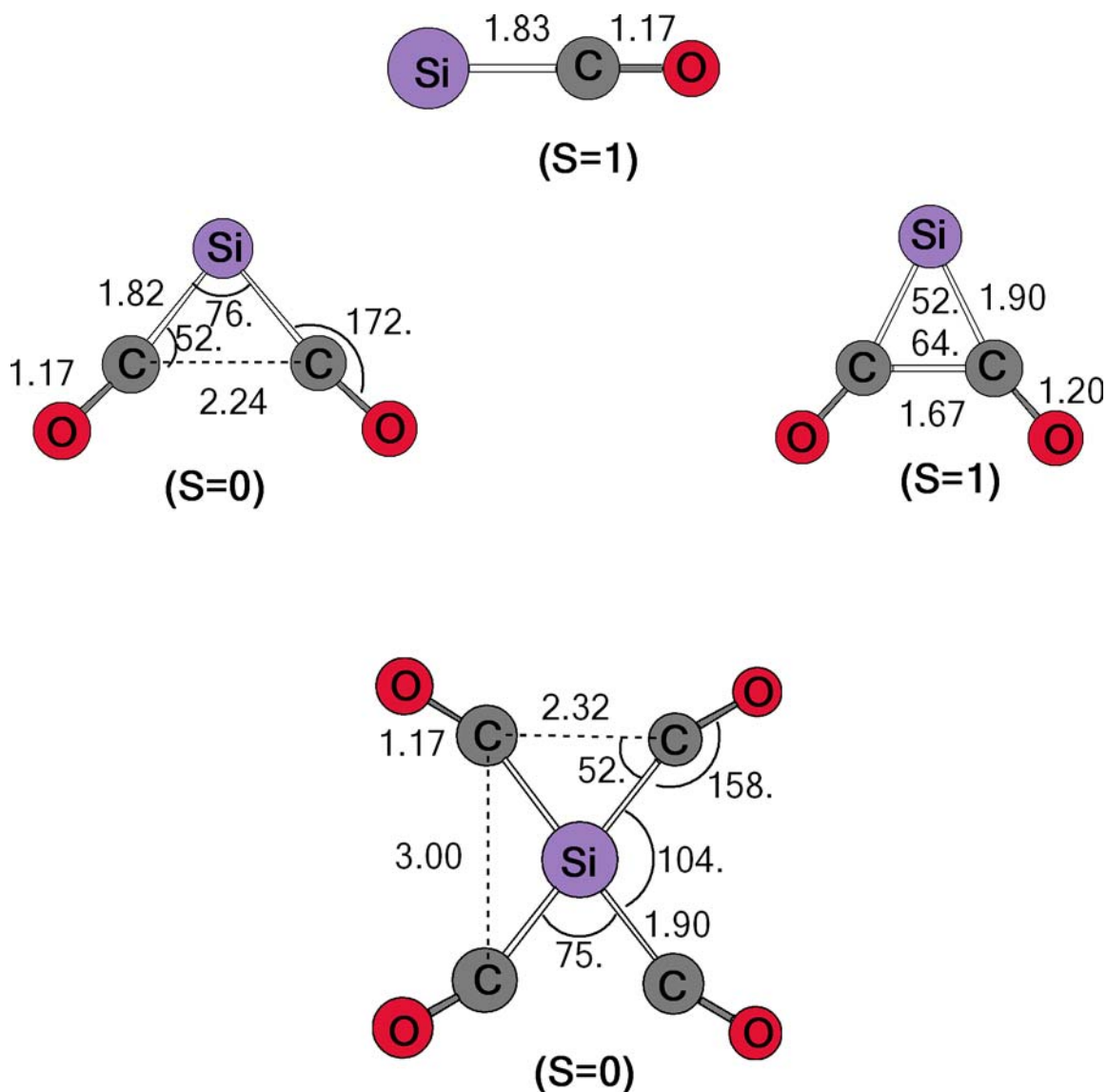
For the search of stable molecules, we considered species involving 1, 2 and 4 CO groups, assigning a priori for each candidate species reasonable symmetries: C<sub>∞v</sub> for Si(CO), C<sub>2v</sub> and D<sub>∞h</sub> for Si(CO)<sub>2</sub>; D<sub>2h</sub>, D<sub>4h</sub>, D<sub>2</sub> and T<sub>d</sub> for Si(CO)<sub>4</sub>. For all of them, we found structures which are more stable than the corresponding reactants.

### 3.1 Geometry and energetics

Table 1 compares the energies  $E$  of the considered geometries with those of the corresponding reactants ( $E$  is measured with respect to the constituting atoms;  $\Delta E$  are the reaction enthalpies of the compounds with respect to reactants), while Fig. 1 shows the corresponding geometries (with internuclear distances in ångstroms and angles in degrees) of the most stable silicon carbonyls Si(CO)<sub>*n*</sub> ( $n = 1, 2, 4$ ).

In order to be considered as molecules, these structures must, however, be stable with respect to small nuclear displacements. We have thus performed a normal mode analysis of the vibrational spectra of all considered molecules. Of them, only Si(CO), Si(CO)<sub>2</sub> in C<sub>2v</sub> symmetry (both singlet and triplet state structures) and Si(CO)<sub>4</sub> in D<sub>2h</sub> symmetry (only singlet state structure) were found to have all normal modes with real frequencies. We have thus focused our attention only on the minimum structures. In Si(CO) the Si–C distance is 1.83 Å, and the C–O bond is 1.17 Å. These bond length values are in good agreement with the optimized Si–C (1.810–1.835 Å) and C–O (1.157–1.161 Å) bond lengths at five theoretical levels (i.e. CASSCF, B3LYP, B3P86, B3PW91, MP2) [8].

In Si(CO)<sub>2</sub> in C<sub>2v</sub> symmetry ( $S = 0$ ) the internuclear distances are the following: 1.82 Å for the Si–C distance, 2.24 Å for the C–C distance, and 1.17 Å for the C–O bond. All these values are in good agreement with those calculated (Si–C: 1.8133 Å; C–O: 1.1500–1.1638 Å) at four theoretical levels



**Fig. 1** The optimized structures of the most stable complexes resulting from the reaction of silicon with 1, 2, and 4 CO ligands. The spin multiplicity of the electronic state is also reported

(i.e. B3LYP, B3P86, B3PW91, MP2) in Ref. [8]. The  $\widehat{\text{CSiC}}$  ( $76^\circ$ ) and  $\widehat{\text{SiCO}}$  ( $172^\circ$ ) angles fall within  $75.39^\circ$ – $79.43^\circ$  and  $170.68^\circ$ – $171.83^\circ$ , respectively, calculated at four theoretical levels in Ref. [8].

In  $\text{Si}(\text{CO})_2$  ( $S = 1$ ) the Si–C bond lengthens to 1.90 Å, while the C–C distance shortens to 1.67 Å; CO bond length is 1.20 Å. As a consequence, the  $\widehat{\text{CSiC}}$  angle decreases to  $52^\circ$ . Analogously, a good comparison can be found with the corresponding calculated data in Ref. [8].

In  $\text{Si}(\text{CO})_4$  in  $D_{2h}$  symmetry ( $S = 0$ ) the Si–C and C–C distances are somewhat larger (1.90 Å and 2.32 Å, respectively) than in  $\text{Si}(\text{CO})_2$  in  $C_{2v}$  symmetry ( $S = 0$ ), while the C–O distance, 1.17 Å, is unchanged. The  $\widehat{\text{CSiC}}$  angle is  $75^\circ$  and the  $\widehat{\text{SiCO}}$  angle is  $158^\circ$ . No theoretical and experimental values for  $\text{Si}(\text{CO})_4$  geometrical parameters have been

reported so far. However, the DFT method, with the Becke88 and Perdew86 functional, and the basis sets used in this work, are reliable and sufficiently accurate, as shown by  $\text{Si}(\text{CO})$  and  $\text{Si}(\text{CO})_2$  results, which are in good agreement with those calculated with different methods [8].

$\text{Si}(\text{CO})_2 S = 1$  is less stable than the corresponding  $S = 0$  by 1.50 eV (about  $35 \text{ kcal mol}^{-1}$ ), as shown in Table 1, and in agreement with Ref. [8] ( $34$ – $42 \text{ kcal mol}^{-1}$  with different theoretical methods). Within calculation accuracy  $\text{Si}(\text{CO})_4$  in  $D_{2h}$  symmetry has the same binding energy as  $\text{Si}(\text{CO})_2$  in  $C_{2v}$  symmetry ( $S = 0$ ) plus 2 CO. From Table 1, the dissociation energy of the  $\text{Si}(\text{CO})$  species in Si plus CO is about  $44 \text{ kcal mol}^{-1}$ , and it is very close to the values from Ref. [8] ( $36.5$ – $41.5 \text{ kcal mol}^{-1}$  from DFT methods,  $42.5 \text{ kcal mol}^{-1}$  from CASSCF-MP2 methods), and slightly overestimated with respect to CCSD(T) value ( $26.1 \text{ kcal mol}^{-1}$ )

**Table 1** Binding energy  $E$  with respect to atoms of the reacting molecules and of the corresponding compounds, and reaction enthalpy  $\Delta E$  of the compounds with respect to reactants

Reactants	$E$ (eV)	Product	Symmetry	$E$ (eV)	$\Delta E$ (eV)	$\Delta E$ (kcal mol <sup>-1</sup> )
1 CO + Si	-15.54	Si(CO) <sup>a</sup>	C <sub>∞v</sub>	-17.45	-1.91	-44.0
2 CO + Si	-30.23	Si(CO) <sub>2</sub>	C <sub>2v</sub>	-33.49	-3.26	-75.2
2 CO + Si	-30.23	Si(CO) <sub>2</sub> <sup>a</sup>	C <sub>2v</sub>	-31.99	-1.76	-40.6
2 CO + Si	-30.23	Si(CO) <sub>2</sub> <sup>a</sup>	D <sub>∞h</sub>	-31.30	-1.07	-24.7
2 CO + Si	-30.23	Si(CO) <sub>2</sub>	D <sub>∞h</sub>	-31.05	-0.82	-18.9
4 CO + Si	-59.62	Si(CO) <sub>4</sub>	D <sub>2h</sub>	-62.85	-3.23	-74.5
4 CO + Si	-59.62	Si(CO) <sub>4</sub>	D <sub>4h</sub>	-62.34	-2.72	-62.7
4 CO + Si	-59.62	Si(CO) <sub>4</sub> <sup>b</sup>	D <sub>2</sub>	-60.29	-0.67	-15.5
4 CO + Si	-59.62	Si(CO) <sub>4</sub> <sup>b</sup>	T <sub>d</sub>	-60.04	-0.42	-9.7

Unless otherwise specified, the product is in singlet state

<sup>a</sup>Triplet state

<sup>b</sup>Quintet state

<sup>c</sup>For the silicon atom in electronic state <sup>3</sup>P, a ground state energy of -0.84 eV has been assumed, as recommended by Baerends et al. [26]

[8]. However, this value indicates that the Si–CO bonding should not be considered a weak interaction.

The first dissociation energy for Si(CO)<sub>2</sub> ( $S = 1$ ) is about 31 kcal mol<sup>-1</sup>, comparable to the values from Ref. [8] (23.6–28.3 kcal mol<sup>-1</sup> from DFT methods, 24.8 from MP2), and smaller than the second CO dissociation energy by about 13 kcal mol<sup>-1</sup>, as in Ref. [8], i.e. the first CO binding energy of an Si atom is greater than the second CO one. At variance with Ref. [8], it is very important to note that the third or fourth CO binding energy of Si is NOT smaller than the second one. In particular, the dissociation energy of Si(CO)<sub>4</sub> into Si(CO) plus three CO is about 30 kcal mol<sup>-1</sup>, i.e. the second CO binding energy of an Si atom is as large as the fourth CO one. From the viewpoint of the average single CO binding energy, the average value (19 kcal mol<sup>-1</sup>) for the Si(CO)<sub>4</sub> species is smaller than that (38 kcal mol<sup>-1</sup>) for the Si(CO)<sub>2</sub> species, the latter being smaller than that (44 kcal mol<sup>-1</sup>) for the Si(CO) species.

## 3.2 Orbital analysis and the nature of the silicon-carbonyl bond

### 3.2.1 Electronic structure

#### Si(CO) : Silicon carbonyl ( $S = 1$ )

This system has a triplet ( $S = 1$ ) ground state (<sup>3</sup>A<sub>2</sub>). The linear molecule has been chosen to lie along the  $z$  axis and descent to C<sub>2v</sub> symmetry has been used to analyze the one-electron molecular orbitals obtained by spin unrestricted calculations in terms of percentage composition based on Mulliken population. The 4a<sub>1</sub> HOMO orbital represents a bonding combination of the doubly occupied 3σ orbital on CO (20%) with 3s (57%) and 3p<sub>z</sub> (20%) orbitals on Si. Above the HOMO, the singly occupied 2b<sub>1</sub> and 2b<sub>2</sub> α molecular orbitals can be described as Si 3p<sub>x</sub> (61%) and 3p<sub>y</sub> (61%), respectively, interacting with CO 2π\* orbitals (36%2π<sub>x</sub>, 36%2π<sub>y</sub>). From the Mulliken gross population analysis, the σ orbital interaction between Si 3s and 3p<sub>z</sub> and CO 3σ causes

a decrease in the 3σ population (from 2.00e to 1.45e) and a corresponding increase in the Si 3p<sub>z</sub> (from 0.00e to 0.60e) Mulliken population. The π interaction is a π-backdonation from Si 3p<sub>x</sub>, 3p<sub>y</sub>, whose Mulliken populations decrease from 1.00ae to 0.63ae, to CO 2π\*, whose population correspondingly increases from 0.00ae to 0.36 + 0.36ae. The resulting Mulliken total charge on silicon is slightly positive, 0.14e, on C is 0.18e and on O atom is -0.32e. The spin density analysis reveals that roughly one unpaired electron is localized on Si atom (1.28e) and the other is shared by the CO group (C 0.45e, O 0.28e).

#### Si(CO)<sub>2</sub> : Silicon dicarbonyl ( $S = 0$ )

Table 2 shows selected one-electron orbitals obtained by spin restricted calculations for the Si(CO)<sub>2</sub> complex in C<sub>2v</sub> symmetry ( $S = 0$ ) (<sup>1</sup>A<sub>1</sub>). The energies and the percentage composition based on Mulliken population analysis is given in terms of atomic Si<sup>+</sup> ion orbitals and (CO)<sub>2</sub><sup>-</sup> orbitals, with the nature of the single “CO” orbitals reported in parentheses. We promoted the fragments to the ionic configurations Si<sup>+</sup> [(3s)<sup>1</sup>(3p<sub>z</sub>)<sup>0</sup>(3p<sub>x</sub>)<sup>0</sup>(3p<sub>y</sub>)<sup>2</sup>] and (CO)<sub>2</sub><sup>-</sup> with one, namely 5a<sub>1</sub> (2π<sub>ip</sub><sup>\*</sup>), of the 2π\* set of orbitals singly occupied. This change of configuration has an advantage in that the Si 3s acts as acceptor orbital for electrons from (CO)<sub>2</sub><sup>-</sup> singly occupied 5a<sub>1</sub> (“CO” 2π\*) and 4a<sub>1</sub> (“CO” 3σ), and the Si 3p<sub>y</sub> orbital can act as donor orbital to the empty (CO)<sub>2</sub><sup>-</sup>2b<sub>2</sub> (“CO” 2π<sub>y</sub><sup>\*</sup>) orbital. The planar molecule has been chosen to lie on the  $xz$  plane.

The 2b<sub>2</sub> HOMO orbital represents a bonding interaction between silicon 3p<sub>y</sub> (46%) and (CO)<sub>2</sub><sup>-</sup>2b<sub>2</sub> (“CO” 2π<sub>y</sub><sup>\*</sup>) orbital, with a small mixing of Si 3d<sub>yz</sub> orbital. Below the HOMO, the 5a<sub>1</sub> MO can be described as silicon 3s (50%) and 3p<sub>z</sub> (25%) interacting with 5a<sub>1</sub> (CO)<sub>2</sub><sup>-</sup> (“CO” 2π\*) and 4a<sub>1</sub> (CO)<sub>2</sub><sup>-</sup> (“CO” 3σ) orbitals. The low-lying 2a<sub>1</sub> orbital is characterized by a Si 3s (22%) and (CO)<sub>2</sub><sup>-</sup>4a<sub>1</sub> (“CO” 3σ) and 2a<sub>1</sub> (“CO” 2σ), with small contribution from Si 3p<sub>z</sub> (5%). Silicon 3p<sub>x</sub>–“CO” 3σ (4b<sub>1</sub>)-1π (3b<sub>1</sub>) bonding interactions are described by 4b<sub>1</sub> and 3b<sub>1</sub> MOs.

**Table 2** One-electron energies and percentage composition (based on Mulliken population analysis per MO) of the complex Si(CO)<sub>2</sub> (*S* = 0) in C<sub>2v</sub> symmetry

Orbital	ε (eV)	Si <sup>+</sup>	(CO) <sub>2</sub> <sup>-1</sup>
2a <sub>1</sub>	-17.23	22% 3s + 5% 3p <sub>z</sub>	37% 2a <sub>1</sub> (2σ) + 23% 4a <sub>1</sub> (3σ)
2b <sub>1</sub>	-14.20	4% 3p <sub>x</sub>	85% 2b <sub>1</sub> (2σ) + 7% 4b <sub>1</sub> (3σ)
3a <sub>1</sub>	-13.12	8% 3s	62% 2a <sub>1</sub> (2σ) + 23% 4a <sub>1</sub> (3σ)
4a <sub>1</sub>	-12.01	3% 3s	86% 3a <sub>1</sub> (1π <sub>ip</sub> ) + 9% 4a <sub>1</sub> (3σ)
3b <sub>1</sub>	-11.83	9% 3p <sub>x</sub>	11% 2b <sub>1</sub> (2σ) + 60% 3b <sub>1</sub> (1π <sub>ip</sub> ) + 18% 4b <sub>1</sub> (3σ)
1b <sub>2</sub>	-11.62	2% 3d <sub>yz</sub>	96% 1b <sub>2</sub> (1π <sub>y</sub> ) + 1% 2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )
1a <sub>2</sub>	-11.29		99% 1a <sub>2</sub> (1π <sub>y</sub> )
4b <sub>1</sub>	-10.47	2% 3d <sub>xz</sub> + 8% 3p <sub>x</sub>	4% 2b <sub>1</sub> (2σ) + 37% 3b <sub>1</sub> (1π <sub>ip</sub> ) + 49% 4b <sub>1</sub> (3σ)
5a <sub>1</sub>	-7.88	50% 3s + 25% 3p <sub>z</sub>	15% 4a <sub>1</sub> (3σ) + 10% 5a <sub>1</sub> (2π <sub>ip</sub> <sup>*</sup> )
2b <sub>2</sub> HOMO	-5.65	1% 3d <sub>yz</sub> + 46% 3p <sub>y</sub>	3% 1b <sub>2</sub> (1π <sub>y</sub> ) + 51% 2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )
6a <sub>1</sub> LUMO	-3.78	4% 3s + 10% 3p <sub>z</sub>	74% 5a <sub>1</sub> (2π <sub>ip</sub> <sup>*</sup> )
2a <sub>2</sub>	-2.49	4% 3d <sub>xy</sub>	94% 2a <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )
5b <sub>1</sub>	-2.17	4% 3d <sub>xz</sub> + 36% 3p <sub>x</sub>	53% 5b <sub>1</sub> (2π <sub>ip</sub> <sup>*</sup> )
3b <sub>2</sub>	-1.33	5% 3d <sub>yz</sub> + 54% 3p <sub>y</sub>	37% 2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )

Selected orbitals involving Si atom in terms of Si<sup>+</sup> and (CO)<sub>2</sub><sup>-1</sup> fragments are reported. The “CO” orbital character is shown in parentheses

**Table 3** Mulliken gross population of valence MO of Si<sup>+</sup> and (CO)<sub>2</sub><sup>-1</sup> fragments in Si(CO)<sub>2</sub> (*S* = 0) in different irreducible representations in C<sub>2v</sub> point group

	A <sub>1</sub>	A <sub>2</sub>	B <sub>1</sub>	B <sub>2</sub>
Si <sup>+</sup>	3s 1.65 3d <sub>z<sup>2</sup></sub> 0.02 3d <sub>x<sup>2</sup>-y<sup>2</sup></sub> 0.03 3p <sub>z</sub> 0.62	3d <sub>xy</sub> 0.01	3d <sub>xz</sub> 0.09 3p <sub>x</sub> 0.43	3d <sub>yz</sub> 0.03 3p <sub>y</sub> 0.96
(CO) <sub>2</sub> <sup>-1</sup>	4a <sub>1</sub> (3σ) 1.54 3a <sub>1</sub> (1π <sub>ip</sub> ) 1.94 5a <sub>1</sub> (2π <sub>ip</sub> <sup>*</sup> ) 0.22	1a <sub>2</sub> (1π <sub>y</sub> ) 1.98	4b <sub>1</sub> (3σ) 1.48 3b <sub>1</sub> (1π <sub>ip</sub> ) 1.96	2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> ) 1.03
Gross charge	Si 0.15	C 0.26	O -0.33	

Table 3 shows the Mulliken gross population of Si<sup>+</sup> and (CO)<sub>2</sub><sup>-1</sup> fragment orbitals in the complex, in different irreducible representations in C<sub>2v</sub> point group.

From Table 3, we see that the σ and π<sub>ip</sub> orbital interactions between Si 3s and 3p<sub>z</sub> and (CO)<sub>2</sub><sup>-1</sup> (“CO” 2π<sup>\*</sup>) and (“CO” 3σ) orbitals in the 5a<sub>1</sub> – 2a<sub>1</sub> MOs causes a decrease in the (CO)<sub>2</sub><sup>-1</sup> 4a<sub>1</sub> (3σ) Mulliken population, which reduces from 2e to 1.54e, and 5a<sub>1</sub> (“CO” 2π<sup>\*</sup>) Mulliken population, which reduces from 1e to 0.22e, and an increase in the Si 3s (from 1e to 1.65e) and 3p<sub>z</sub> (from 0e to 0.62e) Mulliken populations. The Si 3d<sub>z<sup>2</sup></sub> and 3d<sub>x<sup>2</sup>-y<sup>2</sup></sub> get slightly populated by 0.02e and 0.03e, respectively. In B<sub>1</sub> symmetry, the σ orbital interaction between Si 3p<sub>x</sub>, 3d<sub>xz</sub> and (CO)<sub>2</sub><sup>-1</sup> 4b<sub>1</sub> (“CO” 3σ character) causes a decrease in the (CO)<sub>2</sub><sup>-1</sup> 4b<sub>1</sub> Mulliken population which reduces from 2e to 1.48e, and a corresponding increase in the Si 3p<sub>x</sub> (0.43e) and 3d<sub>xz</sub> (0.09e) populations.

Finally, in B<sub>2</sub> symmetry, the π out-of-plane interaction between the doubly occupied Si 3p<sub>y</sub> and (CO)<sub>2</sub><sup>-1</sup> 2b<sub>2</sub> empty orbital (“CO” 2π<sup>\*</sup> character), causes a decrease in the Si 3p<sub>y</sub> population (from 2e to 0.96e) and a corresponding increase in the 2b<sub>2</sub> population (from 0e to 1.03e). The Si 3d<sub>yz</sub> also acquires a small population (0.03e).

In A<sub>2</sub> symmetry very weak interaction between Si 3d<sub>xy</sub> (0.01e) and (CO)<sub>2</sub><sup>-1</sup> 1a<sub>2</sub> (1.98e) can be observed.

We note that silicon 3d<sub>xz</sub> orbital participates in the σ bonding with the two carbonyl set.

The resulting Mulliken total charge on Si atom is slightly positive, 0.15e, on C is 0.26e, and on O atom is -0.33e.

### Si(CO)<sub>2</sub> : c-silicodiketone (*S* = 1)

In Table 4 selected one-electron orbitals (only α-spin) obtained by spin unrestricted calculations of the Si(CO)<sub>2</sub> complex (*S* = 1) (<sup>3</sup>B<sub>2</sub>) are shown. The energies and the percentage composition based on Mulliken population analysis is given in terms of atomic Si [(3s)<sup>2</sup>(3p<sub>x</sub>)<sup>0</sup>(3p<sub>y</sub>)<sup>1</sup>(3p<sub>z</sub>)<sup>1</sup>] and (CO)<sub>2</sub> orbitals.

The 2b<sub>2</sub> molecular orbital is singly occupied and represents a bonding interaction between silicon 3p<sub>y</sub> (44%) and (CO)<sub>2</sub> 2b<sub>2</sub> (2π<sub>y</sub><sup>\*</sup>). The singly occupied 6a<sub>1</sub> molecular orbital can be described as silicon 3p<sub>z</sub> (27%) interacting with (CO)<sub>2</sub> 5a<sub>1</sub> (“CO” 2π<sub>ip</sub><sup>\*</sup>). Silicon 3p<sub>x</sub> – “CO” 3σ (4b<sub>1</sub>) bonding interaction is described by 4b<sub>1</sub> HOMO orbital, while 5a<sub>1</sub> represents a bonding interaction between silicon 3s (52%),

**Table 4** One-electron energies and percentage composition (based on Mulliken population analysis per MO) of the Si(CO)<sub>2</sub> (*S* = 1) complex in C<sub>2v</sub> symmetry

Orbital	ε (eV)	Si	(CO) <sub>2</sub>
2a <sub>1</sub>	-18.99	14% 3s + 5% 3p <sub>z</sub>	70% 2a <sub>1</sub> (2σ) + 6% 4a <sub>1</sub> (3σ)
2b <sub>1</sub>	-13.33	2% 3p <sub>x</sub>	93% 2b <sub>1</sub> (2σ) + 2% 4b <sub>1</sub> (3σ)
3a <sub>1</sub>	-12.93	20% 3s	25% 2a <sub>1</sub> (2σ) + 18% 3a <sub>1</sub> (1π <sub>ip</sub> ) + 36% 4a <sub>1</sub> (3σ)
4a <sub>1</sub>	-12.12	4% 3s	78% 3a <sub>1</sub> (1π <sub>ip</sub> ) + 16% 4a <sub>1</sub> (3σ)
1b <sub>2</sub>	-11.67	3% 3p <sub>y</sub>	95% 1b <sub>2</sub> (1π <sub>y</sub> )
3b <sub>1</sub>	-10.78	8% 3p <sub>x</sub>	75% 3b <sub>1</sub> (1π <sub>ip</sub> ) + 10% 4b <sub>1</sub> (3σ)
1a <sub>2</sub>	-10.65		100% 1a <sub>2</sub> (1π <sub>y</sub> )
5a <sub>1</sub>	-9.21	52% 3s + 14% 3p <sub>z</sub>	29% 4a <sub>1</sub> (3σ) + 4% 5a <sub>1</sub> (2π <sub>ip</sub> <sup>*</sup> )
4b <sub>1</sub> HOMO	-8.31	2% 3d <sub>xz</sub> + 16% 3p <sub>x</sub>	22% 3b <sub>1</sub> (1π <sub>ip</sub> ) + 57% 4b <sub>1</sub> (3σ)
2b <sub>2</sub> SOMO	-6.47	44% 3p <sub>y</sub>	49% 2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )
6a <sub>1</sub> SOMO	-4.86	27% 3p <sub>z</sub>	61% 5a <sub>1</sub> (2π <sub>ip</sub> )
3b <sub>2</sub> LUMO	-2.87	3% 3d <sub>yz</sub> + 56% 3p <sub>y</sub>	45% 2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )
5b <sub>1</sub>	-1.86	7% 3d <sub>xz</sub> + 82% 3p <sub>x</sub>	18% 4b <sub>1</sub> (3σ)
2a <sub>2</sub>	-1.61	4% 3d <sub>xy</sub>	95% 2a <sub>2</sub> (2π <sub>x</sub> <sup>*</sup> )

Selected orbitals (only α-spin) involving Si atom in terms of Si (3s<sup>2</sup>3p<sub>x</sub><sup>0</sup>3p<sub>y</sub><sup>1</sup>3p<sub>z</sub><sup>1</sup>) and (CO)<sub>2</sub> fragments are reported. The “CO” orbital character is shown in parentheses

3p<sub>z</sub> (14%), and (CO)<sub>2</sub>4a<sub>1</sub> (“CO” 3σ) and 5a<sub>1</sub> (2π<sub>ip</sub><sup>\*</sup>). Table 5 shows the Mulliken gross population of Si and (CO)<sub>2</sub> fragment orbitals in Si(CO)<sub>2</sub> (*S* = 1) complex.

In A<sub>1</sub> symmetry we see that the σ and π<sub>ip</sub> orbital interactions between Si 3s and 3p<sub>z</sub> and (CO)<sub>2</sub> 3σ and 2π<sub>ip</sub><sup>\*</sup> orbitals cause a decrease in the (CO)<sub>2</sub> 4a<sub>1</sub> (3σ) Mulliken population, which reduces from 2e to 1.80e, and in both Si 3s (from 2e to 1.81e) and 3p<sub>z</sub> (from 1e to 0.67e) Mulliken populations. The (CO)<sub>2</sub> 5a<sub>1</sub> (“CO” 2π<sub>ip</sub><sup>\*</sup>) gets populated by 0.77e, due to Si 3s and 3p<sub>z</sub> donation. Moreover, a single point calculation performed on the two CO at the same geometry as in the complex has shown that CO (1) 2π<sub>ip</sub><sup>\*</sup> orbital is populated by donation from CO (2) 3σ orbital, and viceversa, thus indicating a bonding interaction between the two CO already without the presence of Si. In B<sub>1</sub> symmetry, the σ orbital interaction between Si 3p<sub>x</sub>, 3d<sub>xz</sub> and (CO)<sub>2</sub> 4b<sub>1</sub> (“CO” 3σ) causes a decrease in the (CO)<sub>2</sub> 4b<sub>1</sub> Mulliken population from 2e to 1.38e and a corresponding increase in the Si 3p<sub>x</sub> (0.54e) and 3d<sub>xz</sub> (0.08e) populations. In B<sub>2</sub> symmetry, the π out-of-plane interaction between the singly occupied Si 3p<sub>y</sub> and (CO)<sub>2</sub> 2b<sub>2</sub> empty orbital (“CO” 2π<sub>y</sub><sup>\*</sup>) causes a decrease in the Si 3p<sub>y</sub> population (from 1e to 0.50e) and a corresponding increase in the 2b<sub>2</sub> population (from 0e to 0.51e). For the Si(CO)<sub>2</sub> complex (*S* = 1), the resulting Mulliken total charge on Si is more positive, 0.29e, on C is 0.23e and on O atom is -0.37e. The spin density is distributed over Si (0.76e), C (0.26e) and O (0.36e) atoms.

In the Si(CO)<sub>2</sub> complex, the nature of the Si bonding is consistent with the use of sp<sup>2</sup>-hybridized orbitals on silicon. For the silicon dicarbonyl complex (*S* = 0) the p orbital at the Si atom is doubly occupied and a π back-donation of lone-pair electrons from this p orbital to the vacant 2π<sup>\*</sup> CO orbitals occurs. In the c-silicodiketone complex (*S* = 1) complex both the p orbital and one of the sp<sup>2</sup> orbital at the Si atom are singly occupied and a π back-donation of the two electrons from Si into the vacant 2π<sup>\*</sup> (in-plane and out-of-plane) CO orbitals takes place.

#### Si(CO)<sub>4</sub>: Silicon tetracarbonyl (*S* = 0)

Table 6 shows selected one-electron orbitals obtained by spin restricted calculations for the Si(CO)<sub>4</sub> complex in D<sub>2h</sub> symmetry (<sup>1</sup>A<sub>1g</sub>). The energies and the percentage composition based on Mulliken population analysis is given in terms of atomic Si<sup>2+</sup> ion orbitals and (CO)<sub>4</sub><sup>2-</sup> orbitals, with the nature of the single “CO” orbitals reported in parentheses. We promoted the fragments to the ionic configurations Si<sup>2+</sup> [(3s)<sup>0</sup>(3p<sub>z</sub>)<sup>0</sup>(3p<sub>x</sub>)<sup>0</sup>(3p<sub>y</sub>)<sup>2</sup>] and (CO)<sub>4</sub><sup>2-</sup> (with one, namely 5a<sub>1g</sub>, of the 2π<sup>\*</sup> set of orbitals doubly occupied). This change of configuration has the advantage that the Si 3s is emptied and acts as acceptor orbital for electrons from (CO)<sub>4</sub><sup>2-</sup> 5a<sub>1g</sub> (“CO” 2π<sup>\*</sup>), and the Si 3p<sub>y</sub> orbital can act as donor orbital to the empty (CO)<sub>4</sub><sup>2-</sup> 2b<sub>2u</sub> (“CO” 2π<sup>\*</sup>) orbital.

The planar molecule has been chosen to lie on the xz plane.

The 5a<sub>1g</sub> HOMO orbital can be described as silicon 3s (22%), with small mixing of Si 3d<sub>z<sup>2</sup></sub> and 3d<sub>x<sup>2</sup>-y<sup>2</sup></sub>, interacting with 5a<sub>1g</sub> (CO)<sub>4</sub><sup>2-</sup> (“CO” 2π<sup>\*</sup>) and 4a<sub>1g</sub> (CO)<sub>4</sub><sup>2-</sup> (“CO” 3σ) orbitals. The low-lying 2a<sub>1g</sub> orbital is characterized by a Si 3s and (CO)<sub>4</sub><sup>2-</sup> 4a<sub>1g</sub> (“CO” 3σ) and 5a<sub>1g</sub> (“CO” 2π<sup>\*</sup>) orbital combination. Below the HOMO, the 2b<sub>2u</sub> MO represents a bonding interaction between silicon 3p<sub>y</sub> (49%) and (CO)<sub>4</sub><sup>2-</sup> 2b<sub>2u</sub> (“CO” 2π<sup>\*</sup>) orbital, while the 4b<sub>2g</sub> is mainly (CO)<sub>4</sub><sup>2-</sup> 4b<sub>2g</sub> (“CO” 3σ) orbital (79%) mixing with silicon 3d<sub>xz</sub> (9%). Silicon 3p<sub>x</sub> – “CO” 3σ–2σ bonding interactions are described by orbitals 3b<sub>3u</sub> and 2b<sub>3u</sub>, while the corresponding silicon 3p<sub>z</sub> – “CO” 3σ – 2σ bonding interactions are represented by 3b<sub>1u</sub> and 2b<sub>1u</sub> MOs.

Table 7 shows the Mulliken gross population of Si<sup>2+</sup> and (CO)<sub>4</sub><sup>2-</sup> fragment orbitals in the silicon tetracarbonyl complex, in different irreducible representations in D<sub>2h</sub> point group.

From Table 7 we see that the π<sub>i.p.</sub> orbital interaction in the HOMO 5a<sub>1g</sub> causes a decrease in the (CO)<sub>4</sub><sup>2-</sup> 5a<sub>1g</sub> Mulliken

**Table 5** Mulliken gross population of valence MO of Si and (CO)<sub>2</sub> fragments in Si(CO)<sub>2</sub> (*S* = 1) in different irreducible representations in C<sub>2v</sub> point group

	A <sub>1</sub>	A <sub>2</sub>	B <sub>1</sub>	B <sub>2</sub>
Si	3s 1.81 3d <sub>z<sup>2</sup></sub> 0.04 3d <sub>x<sup>2</sup>-y<sup>2</sup></sub> 0.02 3p <sub>z</sub> 0.67	3d <sub>xy</sub> 0.00	3d <sub>xz</sub> 0.08 3p <sub>x</sub> 0.54	3d <sub>yz</sub> 0.03 3p <sub>y</sub> 0.50
(CO) <sub>2</sub>	4a <sub>1</sub> (3σ)1.80 3a <sub>1</sub> (1π <sub>ip</sub> )1.99 5a <sub>1</sub> (2π <sub>ip</sub> <sup>*</sup> )0.77	1a <sub>2</sub> (1π <sub>y</sub> )1.99	4b <sub>1</sub> (3σ)1.38 3b <sub>1</sub> (1π <sub>ip</sub> )1.96	2b <sub>2</sub> (2π <sub>y</sub> <sup>*</sup> )0.51
Gross charge	Si 0.29	C 0.23	O -0.37	
Spin density	Si 0.76	C 0.26	O 0.36	

**Table 6** One-electron energies and percentage composition (based on Mulliken population analysis per MO) of the Si(CO)<sub>4</sub> (*S* = 0) complex

Orbital	ε (eV)	Si <sup>2+</sup>	(CO) <sub>4</sub> <sup>2-</sup>
2a <sub>1g</sub>	-17.64	30% 3s	28%2a <sub>1g</sub> (2σ) + 38%4a <sub>1g</sub> (3σ) + 1%5a <sub>1g</sub> (2π <sup>*</sup> )
2b <sub>1u</sub>	-15.72	18% 3p <sub>z</sub>	51%2b <sub>1u</sub> (2σ) + 26%4b <sub>1u</sub> (3σ) + 3%3b <sub>1u</sub> (1π)
2b <sub>3u</sub>	-14.54	10% 3p <sub>x</sub>	76%2b <sub>3u</sub> (2σ) + 9%4b <sub>3u</sub> (3σ) + 4%3b <sub>3u</sub> (1π)
2b <sub>2g</sub>	-14.07	1% 3d <sub>xz</sub>	95%2b <sub>2g</sub> (2σ) + 3%4b <sub>2g</sub> (3σ)
3a <sub>1g</sub>	-13.32	7% 3s	72%2a <sub>1g</sub> (2σ) + 20%4a <sub>1g</sub> (3σ)
3b <sub>1u</sub>	-12.81	9% 3p <sub>z</sub>	34%4b <sub>1u</sub> (3σ) + 47%2b <sub>1u</sub> (2σ) + 9%3b <sub>1u</sub> (1π)
3b <sub>3u</sub>	-12.34	13% 3p <sub>x</sub>	52%3b <sub>3u</sub> (1π) + 19%2b <sub>3u</sub> (2σ) + 15%4b <sub>3u</sub> (3σ)
4a <sub>1g</sub>	-11.99	1% 3d <sub>x<sup>2</sup>-y<sup>2</sup></sub>	97%3a <sub>1g</sub> (1π)
1b <sub>2u</sub>	-11.86	4% 3p <sub>y</sub>	94%1b <sub>2u</sub> (1π) + 1%2b <sub>2u</sub> (2π <sup>*</sup> )
1b <sub>3g</sub>	-11.69	1% 3d <sub>yz</sub>	99%1b <sub>3g</sub> (1π)
4b <sub>1u</sub>	-11.60	1% 3p <sub>z</sub>	87%3b <sub>1u</sub> (1π) + 10%4b <sub>1u</sub> (3σ) + 1%2b <sub>1u</sub> (2σ)
1b <sub>1g</sub>	-11.56	1% 3d <sub>xy</sub>	99%1b <sub>1g</sub> (1π)
1a <sub>1u</sub>	-11.49		99%1a <sub>1u</sub> (1π)
3b <sub>2g</sub>	-11.42	1% 3d <sub>xz</sub>	93%3b <sub>2g</sub> (1π) + 4%4b <sub>2g</sub> (3σ) + 1%2b <sub>2g</sub> (2σ)
4b <sub>3u</sub>	-10.51	8% 3p <sub>x</sub>	45%4b <sub>3u</sub> (3σ) + 41%3b <sub>3u</sub> (1π) + 4%2b <sub>3u</sub> (2σ)
4b <sub>2g</sub>	-9.53	9% 3d <sub>xz</sub>	79%4b <sub>2g</sub> (3σ) + 6%3b <sub>2g</sub> (1π) + 3%2b <sub>2g</sub> (2σ)
2b <sub>2u</sub>	-6.35	49% 3p <sub>y</sub>	45%2b <sub>2u</sub> (2π <sup>*</sup> ) + 5%1b <sub>2u</sub> (1π)
5a <sub>1g</sub> HOMO	-5.94	22% 3s + 3%3d <sub>z<sup>2</sup></sub> + 2%3d <sub>x<sup>2</sup>-y<sup>2</sup></sub>	54%5a <sub>1g</sub> (2π <sup>*</sup> ) + 22%4a <sub>1g</sub> (3σ)
2b <sub>3g</sub> LUMO	-4.00	7% 3d <sub>yz</sub>	92%2b <sub>3g</sub> (2π <sup>*</sup> )
5b <sub>1u</sub>	-3.45	4% 3p <sub>z</sub>	83%5b <sub>1u</sub> (2π <sup>*</sup> ) + 9%4b <sub>1u</sub> (3σ)
2b <sub>1g</sub>	-2.93	6% 3d <sub>xy</sub>	93%2b <sub>1g</sub> (2π <sup>*</sup> )

Selected orbitals involving Si atom in terms of Si<sup>2+</sup> and (CO)<sub>4</sub><sup>2-</sup> fragments are reported. The “CO” orbital character is shown in parentheses

**Table 7** Mulliken gross population of the most important valence MO of Si<sup>2+</sup> and (CO)<sub>4</sub><sup>2-</sup> fragments in Si(CO)<sub>4</sub> (*S* = 0) in the irreducible representations in D<sub>2h</sub> point group

	A <sub>1g</sub>	B <sub>2g</sub>	B <sub>1u</sub>	B <sub>2u</sub>	B <sub>3u</sub>
Si <sup>2+</sup>	3s 1.18 3d <sub>x<sup>2</sup>-y<sup>2</sup></sub> 0.07 3d <sub>z<sup>2</sup></sub> 0.09	3d <sub>xz</sub> 0.24	3p <sub>z</sub> 0.55	3p <sub>y</sub> 1.06	3p <sub>x</sub> 0.63
(CO) <sub>4</sub> <sup>2-</sup>	4a <sub>1g</sub> (3σ)1.62 5a <sub>1g</sub> (2π <sup>*</sup> )1.13	4b <sub>2g</sub> (3σ)1.72	4b <sub>1u</sub> (3σ)1.41	2b <sub>2u</sub> (2π <sup>*</sup> )0.91	3b <sub>3u</sub> (1π)1.93 4b <sub>3u</sub> (3σ)1.37
Gross charge	Si 0.05	C 0.32	O -0.33		

population which reduces from 2e to 1.13e, and an increase in the Si 3d<sub>z<sup>2</sup></sub> and 3d<sub>x<sup>2</sup>-y<sup>2</sup></sub> population (0.09e and 0.07e, respectively). The remaining (CO)<sub>4</sub><sup>2-</sup> 5a<sub>1g</sub> decrease in Mulliken

population is due to a σ donation to Si 3s, which populates as 1.18e, donations partly coming from (CO)<sub>4</sub><sup>2-</sup> 4a<sub>1g</sub> also, whose population changes from 2e to 1.62e.

In B<sub>2g</sub> symmetry, the  $\sigma$  orbital interaction in 4b<sub>2g</sub> MO causes a decrease in the (CO)<sub>4</sub><sup>2-</sup> 4b<sub>2g</sub> (“CO” 3 $\sigma$  character) Mulliken population which reduces from 2e to 1.72e, and a corresponding increase in the Si 3d<sub>xz</sub> population (0.24e).

In B<sub>1u</sub> symmetry, the Si 3p<sub>z</sub> – “CO” 3 $\sigma$  orbital interaction populates the Si 3p<sub>z</sub> by 0.55e, coming from (CO)<sub>4</sub><sup>2-</sup> 4b<sub>1u</sub> (1.41e) orbital.

In B<sub>3u</sub> symmetry, a similar Si 3p<sub>x</sub> – “CO” 3 $\sigma$  orbital interaction populates the Si 3p<sub>x</sub> by 0.63e, coming from (CO)<sub>4</sub><sup>2-</sup> 3b<sub>3u</sub> (1.93e) and 4b<sub>3u</sub> (from 2e to 1.37e) orbitals.

Finally, the  $\pi$  out-of-plane interaction between the doubly occupied Si 3p<sub>y</sub> and (CO)<sub>4</sub><sup>2-</sup> 2b<sub>2u</sub> empty orbital (“CO” 2 $\pi^*$  character), causes a decrease in the Si 3p<sub>y</sub> population (from 2e to 1.06e) and a corresponding increase in the 2b<sub>2u</sub> population (from 0e to 0.91e).

In this complex silicon, 3d<sub>xz</sub> orbital participates quite largely to the  $\sigma$  bonding with the four carbonyl set.

The resulting Mulliken total charge on Si atom is close to zero (0.05e), on C is 0.32e, and on O atom is –0.33e.

Both the pseudo-square planar structure about Si and the Si bonding in Si(CO)<sub>4</sub> are consistent with the use of sp<sup>2</sup>d-hybridized orbitals on silicon. The donation of lone-pair electrons from CO into the vacant 3d orbitals of silicon is responsible for the sp<sup>2</sup>d hybridization preference over the usual sp<sup>3</sup> one. This complex represents a clear example of d orbital involvement in the formation of a four-coordinated silicon.

### 3.2.2 Bonding energy analysis

The charge rearrangements are a qualitative indication of the bonding interactions, but not a quantitative measure of the corresponding energies. Those are explicitly calculated by the energy decomposition scheme discussed in Sect. 2 and displayed in Table 8 for silicon dicarbonyl and silicon tetracarbonyl complexes.

As already noted above, we use an ionic Si<sup>+</sup> fragment in the valence state 3s<sup>1</sup> 3p <sub>$\sigma$</sub> <sup>0</sup> 3p <sub>$\pi_i$</sub> <sup>0</sup> 3p <sub>$\pi_p$</sub> <sup>2</sup> as found in the Si(CO)<sub>2</sub> situation, and an ionic Si<sup>2+</sup> fragment in the valence state 3s<sup>0</sup> 3p <sub>$\sigma$</sub> <sup>0</sup> 3p <sub>$\pi_i$</sub> <sup>0</sup> 3p <sub>$\pi_p$</sub> <sup>2</sup> as found in the Si(CO)<sub>4</sub> molecule. For ligands, we use (CO)<sub>2</sub><sup>-</sup> as fragment, singly occupying the 5a<sub>1</sub> molecular orbital, and (CO)<sub>4</sub><sup>2-</sup> as fragment, doubly occupying 5a<sub>1g</sub> molecular orbital, and the geometry of the fragments as it occurs in the complexes.

As shown in Table 8, the steric interaction energy  $\Delta E^0$  is strongly attractive for Si(CO)<sub>4</sub> complex, due to the fact that the stabilizing contribution arising from the large attractive interaction between the charged fragments,  $\Delta E_{\text{elstat}}$ , overcomes the positive (destabilizing) Pauli repulsion term,  $\Delta E_{\text{Pauli}}$ . The electrostatic interaction in Si(CO)<sub>2</sub> is much less attractive than in Si(CO)<sub>4</sub>, due to the less charged interacting fragments, and slightly overcomes the positive (destabilizing) Pauli repulsion term, thus causing a slightly attractive steric interaction energy  $\Delta E^0$ .

**Table 8** Decomposition of the bonding energy for the formation of the silicon dicarbonyl (*S* = 0) complex in C<sub>2v</sub> symmetry in terms of Si<sup>1+</sup> and (CO)<sub>2</sub><sup>1-</sup> fragments and of the silicon tetracarbonyl (*S* = 0) complex in D<sub>2h</sub> symmetry in terms of Si<sup>2+</sup> and (CO)<sub>4</sub><sup>2-</sup> fragments

	Si(CO) <sub>2</sub>		Si(CO) <sub>4</sub>
$\Delta E_{\text{Pauli}}$	12.13	$\Delta E_{\text{Pauli}}$	5.46
$\Delta E_{\text{elstat}}$	–13.16	$\Delta E_{\text{elstat}}$	–30.08
$\Delta E^0$	–1.01	$\Delta E^0$	–24.97
$\Delta E_{A_1}$	–13.76	$\Delta E_{A_1g}$	–8.99
$\Delta E_{A_2}$	–0.27	$\Delta E_{B_1g}$	–0.45
$\Delta E_{B_1}$	–3.23	$\Delta E_{B_2g}$	–3.03
$\Delta E_{B_2}$	–2.58	$\Delta E_{B_3g}$	–0.57
		$\Delta E_{A_1u}$	–0.37
		$\Delta E_{B_1u}$	–4.40
		$\Delta E_{B_2u}$	–2.56
		$\Delta E_{B_3u}$	–4.27
$\Delta E_{\text{o.i.}}$	–19.84	$\Delta E_{\text{o.i.}}$	–24.63
$\Delta E_{\text{total}}$	–20.85	$\Delta E_{\text{total}}$	–49.60
$\Delta E_{\text{prep}}$	17.58	$\Delta E_{\text{prep}}$	46.38
$\Delta E$	–3.27	$\Delta E$	–3.22

$\Delta E^0$  is the steric repulsion, the  $\Delta E(\Gamma)$  is the contribution due to orbital interaction in different irreducible representations,  $\Delta E_{\text{o.i.}}$  is the total orbital interaction contribution,  $\Delta E_{\text{total}}$  is the sum of  $\Delta E^0$  and  $\Delta E_{\text{o.i.}}$ . Preparation energies ( $\Delta E_{\text{prep}}$ ) of the fragments and bonding energies ( $\Delta E$ ) of the adducts are also given

From the data reported in Table 8, we note that the  $\Delta E_{A_1}$  term for Si(CO)<sub>2</sub> and the  $\Delta E_{A_1g}$  and  $\Delta E_{B_1u}$  terms for Si(CO)<sub>4</sub> account for  $\sigma$ - and  $\pi_{ip}$ -donation to silicon orbitals. For Si(CO)<sub>2</sub> complex, the donation to the 3s (0.65e), 3p<sub>z</sub> (0.62e), 3d<sub>z<sup>2</sup></sub> (0.02e) and 3d<sub>x<sup>2</sup>–y<sup>2</sup></sub> (0.03e) orbitals gives rise to a large  $\Delta E_{A_1}$  energy contribution (–13.76eV). For Si(CO)<sub>4</sub> complex, the donation to the silicon 3p<sub>z</sub> (0.55e) is responsible for a  $\Delta E_{B_1u}$  term of –4.40 eV, while the donation to the 3s (1.18e), 3d<sub>z<sup>2</sup></sub> (0.09e) and 3d<sub>x<sup>2</sup>–y<sup>2</sup></sub> (0.07e) orbitals gives a –8.99 eV contribution ( $\Delta E_{A_1g}$ ). The energy contributions due to the  $\sigma$ - and  $\pi_{ip}$ -donation to silicon orbitals are, therefore, comparable for the two complexes.

The  $\Delta E_{B_1}$  term for Si(CO)<sub>2</sub> and the  $\Delta E_{B_3u}$  and  $\Delta E_{B_2g}$  terms for Si(CO)<sub>4</sub> are also found to give quite a relevant contribution to the  $\sigma$  bond (–3.23 eV for silicon dicarbonyl and –4.27 and –3.03 eV, respectively, for silicon tetracarbonyl), reflecting the large charge transfers into 3p<sub>x</sub> (0.43e for Si(CO)<sub>2</sub> and 0.63e for Si(CO)<sub>4</sub>), by carbonyl orbitals. For silicon tetracarbonyl complex, more relevant charge transfer occurs into 3d<sub>xz</sub> (0.24e), which accounts for –3.03 eV ( $\Delta E_{B_2g}$ ) energy contribution.

As for the energy terms which account for the  $\pi$  out of plane bond, the  $\Delta E_{B_2}$  for Si(CO)<sub>2</sub> and  $\Delta E_{B_2u}$  plus  $\Delta E_{B_3g}$  for Si(CO)<sub>4</sub> terms are large. As inferred by the charge transfers that occur in these symmetries (1.03e into 2b<sub>2</sub> for Si(CO)<sub>2</sub> and 0.91e into 2b<sub>2u</sub> for Si(CO)<sub>4</sub>), the  $\Delta E_{B_2}$  (–2.58 eV) is comparable to  $\Delta E_{B_2u}$  (–2.56 eV). However,  $\Delta E_{B_2}$  also contains contributions due to small donations to 3d<sub>yz</sub>, while this is measured by  $\Delta E_{B_3g}$  term (–0.57 eV) in Si(CO)<sub>4</sub> complex.



The  $\Delta E_{A_2}$  term for  $\text{Si}(\text{CO})_2$  and the corresponding  $\Delta E_{B_{1g}}$  term for  $\text{Si}(\text{CO})_4$  are small, due to the fact that these terms contain the contribution due to interaction of  $3d_{xy}$  orbitals of Si with “CO”  $1\pi$  orbitals of the carbonyl frameworks. Only a very small charge transfer into  $3d_{xy}$  is calculated ( $0.01e$  for  $\text{Si}(\text{CO})_2$  and  $0.02e$  for  $\text{Si}(\text{CO})_4$ ).

As a result of all these contributions, the orbital interaction term  $\Delta E_{o.i.}$  is larger than the steric interaction term  $\Delta E^0$  for  $\text{Si}(\text{CO})_2$  ( $-19.84$  vs.  $-1.01$  eV), while the orbital interaction term is smaller than the steric interaction term for  $\text{Si}(\text{CO})_4$  ( $-24.63$  vs.  $-24.97$  eV). However, it is interesting to note that  $\Delta E_{o.i.}$  for  $\text{Si}(\text{CO})_4$  is larger than  $\Delta E_{o.i.}$  for  $\text{Si}(\text{CO})_2$  by  $4.79$  eV, mainly due to the larger charge transfer into Si  $3d_{xz}$  ( $0.24e$ ) and  $3p_x$  ( $0.63e$ ) orbitals in  $\text{Si}(\text{CO})_4$  than in  $\text{Si}(\text{CO})_2$  complex.

In order to calculate the reaction enthalpy  $\Delta E$  for the formation of the complexes, the preparation energy of the fragments has to be taken into account. The preparation energy,  $\Delta E_{\text{prep}}$ , is largely dominated by the energy necessary to excite the silicon atom from the ground electronic configuration to empty or singly occupy the  $3s$  orbital and to doubly populate the  $3p_\pi$  orbital as we calculate in the converged complexes. The remaining value represents the geometry and valence changes of  $(\text{CO})_2^- / (\text{CO})_4^{2-}$ , with one of the  $2\pi^*$  singly or doubly occupied, bonding to Si in the complex. The above detailed analysis of the different contributions to  $\Delta E_{o.i.}$  points out that in the two complexes: (1) the  $\sigma$  interactions between the  $\text{Si}^+ / \text{Si}^{2+}$  and the  $(\text{CO})_2^- / (\text{CO})_4^{2-}$  framework, due to electron donation from carbonyl to silicon are by far dominant and account for most of the bond strength; (2)  $\pi$  out of plane back-donation from silicon  $3p_\pi$  to empty  $\pi^*$  antibonding  $(\text{CO})_4^{2-}$  and  $(\text{CO})_2^-$  orbitals gives a 13% contribution for  $\text{Si}(\text{CO})_2$  and 10% contribution for  $\text{Si}(\text{CO})_4$  to the silicon-carbonyl bond.

Finally, although in  $\text{Si}(\text{CO})_4$  complex, the Mulliken charge on Si is close to zero, due to a net charge donation from  $(\text{CO})_4^{2-}$  to  $\text{Si}^{2+}$ , a not negligible positive charge can be calculated on Si in  $\text{Si}(\text{CO})_2$  complex, i.e. 0.15.

## 4 Discussion

Even though to a certain extent the comments on the results given in the previous part may be considered as an adequate interpretation of the outcomes of calculations, we, however, believe that the ultimate understanding is achieved when calculations are interpreted in terms of characters of general chemistry like nature (covalent, dative or ionic) and multiplicity of bond, and charge distribution on atoms. To this purpose, we focus our attention on geometry, interatomic-distance distribution, and stability.

We shall consider the four candidate species  $\text{Si}(\text{CO})$  ( $S = 1$ ),  $\text{Si}(\text{CO})_2$  ( $S = 0$ ),  $\text{Si}(\text{CO})_2$  ( $S = 1$ ) and  $\text{Si}(\text{CO})_4$ .

The calculated bond distances are summarized in Table 9, while reference bonding enthalpies (in eV) are reported in Table 10.

**Table 9** Calculated bond distances in the considered silicon carbonyls

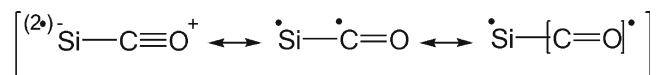
Interatomic pair	Si(CO)	Species Si(CO) <sub>2</sub>		Si(CO) <sub>4</sub>
		S = 0	S = 1	
Si-C	1.83	1.82	1.90	1.90
C-O	1.17	1.17	1.20	1.17
C-C		2.24	1.67	2.32

**Table 10** Some relevant reference bonding enthalpies and bondlengths

Bond	$E_b$ (eV)	Bond length (Å)
C-C	3.59	1.54
C...C	5.25 (in benzene)	1.40
C = C	6.36	1.34
C $\equiv$ C	8.70	1.20
C-O	3.71	1.43
C = O	8.33	1.20
$^-C \equiv O^+$	11.15	1.13
Si-C	3.18	1.87
S-C	2.82	1.82
O=O	5.16	1.21

### 4.1 Si(CO) ( $S = 1$ ): silicon carbonyl

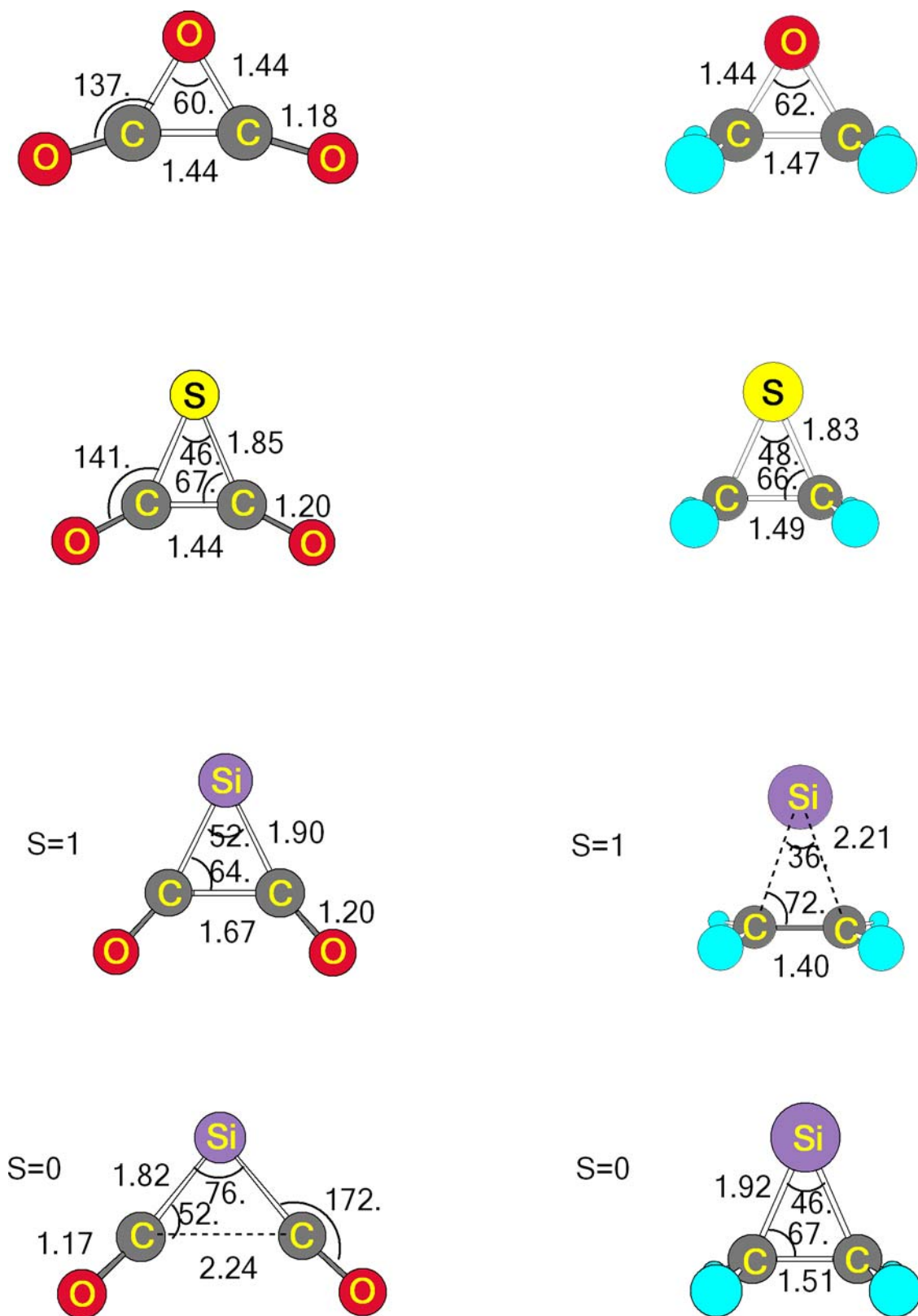
The Si-C distance in  $\text{Si}(\text{CO})$  clearly denotes that an Si-C bond is actually formed. Since the species is in triplet state, it contains two unpaired electrons. Two extreme situations may be hypothesized: either both electrons are on silicon (in species  $(^{2\bullet})^-\text{Si}-C \equiv O^+$ ) or are distributed on carbon and silicon (in species  $\bullet\text{Si}-\bullet\text{C} = \text{O}$  or possibly  $\bullet\text{Si}-[\text{C} = \text{O}]^\bullet$ ). Mulliken charge analysis, giving a spin density of 1.28 on Si, 0.45 on C and 0.28 on O, suggests that  $\text{Si}(\text{CO})$  may be viewed as a kind of resonant structure



The C-O distance, intermediate between that in  $^-C \equiv O^+$  (1.13 Å) and that in aldehydes (1.20 Å), suggests the correctness of this attribution.

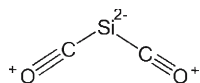
### 4.2 Si(CO)<sub>2</sub>: ( $S = 0$ ) silicon dicarbonyl

The first structure with two CO upon which we focus our attention is  $\text{Si}(\text{CO})_2$  ( $S = 0$ ). The clue for attributing a Lewis formula to this species is the observation that the C-O distance therein (1.17 Å), though larger than in  $^-C \equiv O^+$ , remains, however, shorter than that characteristic of the carbonyl group (1.20 Å in aldehydes). This might be explained assuming that the C-O distance is relaxed with respect to that in  $^-C \equiv O^+$  because the electrostatic reinforcement to

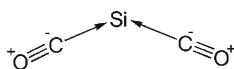


**Fig. 2** The optimized structures of the  $\text{C}_2\text{H}_4\text{X}$  ( $\text{X} = \text{O}, \text{Si}, \text{S}$ ) and  $(\text{CO})_2\text{X}$  ( $\text{X} = \text{O}, \text{Si}, \text{S}$ ) molecules. The spin multiplicity of the electronic state is also reported for silicon based molecules

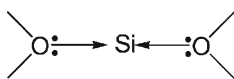
this bond (due to the slight negative net charge on carbon and positive on oxygen) is suppressed because of electronic shift to the silicon atom and the formation of a bond between silicon and carbon because of the unpaired electron on them. The resulting Si–C bond has no strain and is thus characterized by the same internuclear separation characteristic as for the Si–C bond. This compound admits a Lewis formula



which clarifies that the negative charge, originally on carbon in  ${}^{-}\text{C} \equiv \text{O}^{+}$ , has shifted to silicon thus reducing the electrostatic reinforcement to the  $\text{C} \equiv \text{O}$  bonds. In this way, compound  $\text{Si}(\text{CO})_2$  ( $S = 0$ ) may be regarded as *silicon dicarbonyl*. An alternative, but substantially equivalent, description of  $\text{Si}(\text{CO})_2$  ( $S = 0$ ) is in terms of formation of dative bonds via lone pair donation from carbon to silicon:



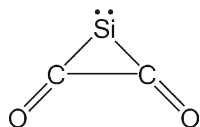
This picture is substantially the same as that advocated for the formation of oxo–oxygen–silicon adducts in siloxanic networks [3–5]:



The stability of silicon dicarbonyl with respect to dissociation  $\text{Si}(\text{CO})_2 \rightarrow \text{Si} + 2\text{CO}$  is essentially conferred by the dative bond. This implies that the bond dissociation energy of the  ${}^{+}\text{O} \equiv \text{C}^{-} \rightarrow \text{Si}$  dative bond is 1.63 eV.

#### 4.3 $\text{Si}(\text{CO})_2$ ( $S = 1$ ): *c*-silicodiketone

Totally different is the internuclear distance distribution in  $\text{Si}(\text{CO})_2$  ( $S = 1$ ). Taking  $\text{Si}(\text{CO})_2$  ( $S = 0$ ) as reference, in the triplet compound the C–O distance coincides with the carbonyl bond length, the C–C distance is appreciably shortened (actually, it is so short as to make not unreasonable the formation of a C–C bond), and the Si–C distance is slightly larger than the Si–C bond length. This situation suggests the following structural formula



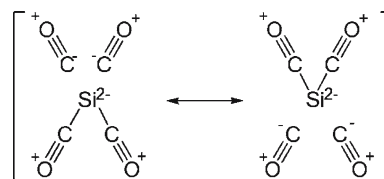
which attributes the increase of the Si–C and C–C distance with respect to the corresponding bond lengths to the strain of the C–Si–C ring. An additional indication of the correctness of this formula is given by the direction of the C–O bond, which points approximately along the bisector of the  $\widehat{\text{Si-C-C}}$  angle, thus suggesting  $sp^2$  hybridization for carbon. This species, whose bare formula might be more properly written as  $\text{Si}[\text{C}_2\text{O}_2]$ , might be called *c-silicodiketone*.

*Mutatis mutandis*, *c*-silicodiketone is indeed the structural analogue of the silacyclopropylidene species studied in refs. [6,7]. Figure 2 shows the structure of  $\text{X}[\text{C}_2\text{H}_4]$  molecules ( $\text{X} = \text{O}$ , oxirane;  $\text{X} = \text{S}$ , thiirane;  $\text{X} = \text{Si}$ , ‘silirane’ or silacyclopropylidene of Ref. [6]), and Table 11 gives their dissociation energies with respect to  $\text{X} + \text{C}_2\text{H}_4$  (the calculations having been carried out at the same level as that adopted in this work).

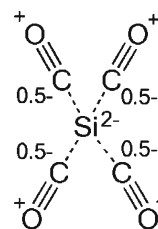
Interestingly enough, the corresponding molecules  $\text{X}[\text{C}_2\text{O}_2]$  (‘*c*Xdiketones’) have well defined minima, that allows them to be considered as chemical species. Of course  $\text{X}[\text{C}_2\text{O}_2]$  ( $\text{X} = \text{O}, \text{S}$ ) are unstable and dissociate spontaneously with  $\text{CO} + \text{XCO}$ ; however,  $\text{Si}[\text{C}_2\text{O}_2]$  is stable with respect to this dissociation, thus suggesting the hypothesis of its possible preparation.

#### 4.4 $\text{Si}(\text{CO})_4$ ( $S = 0$ ): silicon tetracarbonyl

The C–O distance in  $\text{Si}(\text{CO})_4$  indicates that carbon and oxygen are bonded in carbonyl configuration. The distribution of the C–C and Si–C interatomic distances [closer to, but somewhat larger than, the corresponding distances in  $\text{Si}(\text{CO})_2$ ], the binding energy [very close to that of  $\text{Si}(\text{CO})_2 + 2\text{CO}$ ], and the Mulliken charge on silicon [near that on silicon in  $\text{Si}(\text{CO})_2$ ], all together suggest that  $\text{Si}(\text{CO})_4$ , *silicon tetracarbonyl*, may be seen as a resonance structure between the limiting configurations



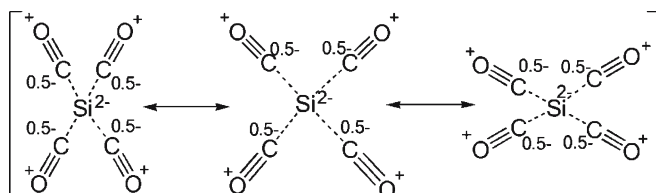
so that it can be described with the following half-bond formula



**Table 11** Binding energy  $E$  with respect to atoms of the reacting molecules and of the corresponding complexes (in  $C_{2v}$  symmetry), and reaction enthalpy  $\Delta E$  of the molecules with respect to reactants

Reactants	$E$ (eV)	Product	$E$ (eV)	$\Delta E$ (eV)
C <sub>2</sub> H <sub>4</sub> + O	-33.43	O(CH <sub>2</sub> ) <sub>2</sub>	-37.54	-4.11
C <sub>2</sub> H <sub>4</sub> + S	-32.49	S(CH <sub>2</sub> ) <sub>2</sub>	-35.24	-2.75
C <sub>2</sub> H <sub>4</sub> + Si	-32.25	Si(CH <sub>2</sub> ) <sub>2</sub> ( $S = 0$ )	-34.45	-2.10
C <sub>2</sub> H <sub>4</sub> + Si	-32.25	Si(CH <sub>2</sub> ) <sub>2</sub> ( $S = 1$ )	-33.40	-1.15
2 CO + O	-31.31	O(CO) <sub>2</sub>	-35.85	-4.54
2 CO + S	-30.37	S(CO) <sub>2</sub>	-32.87	-2.50
2 CO + Si	-30.23	Si(CO) <sub>2</sub> ( $S = 0$ )	-33.49	-3.26
2 CO + Si	-30.23	Si(CO) <sub>2</sub> ( $S = 1$ )	-31.99	-1.76

Even though the square planar Si(CO)<sub>4</sub> complex may actually be viewed as the activated complex allowing the transition from resonance I to resonance II,



actually the same transition is also achieved via a path involving appreciably lower energy-rotation.

The existence of planar tetracoordinated silicon is not so absurd. In fact, X-ray crystallographic investigation on the orthosilicic acid ester led Meyer and Nagorsen [27] to establish it as the first compound with planar tetracoordinated silicon. Later, Schomburg [28] reported evidence of strong distortion of the tetrahedral geometry in a spiro-silicate, bis(tetramethylethylenedioxy)silane. At last, silicon has recently been reported to be a better candidate than carbon for the detection and observation of planar tetracoordinated structures [29].

## 5 Conclusions

Density functional calculations have been performed on a few relatively stable compounds of silicon with CO: Si(CO), Si(CO)<sub>2</sub>, Si(CO)<sub>4</sub> and Si[C<sub>2</sub>O<sub>2</sub>]. For Si(CO)<sub>n</sub> ( $n = 1, 2$ ), our results are in good agreement with the ones in literature, confirming the accuracy of our approach. The thermodynamic stability of a planar tetracoordinated Si(CO)<sub>4</sub> complex has been shown. To the best of our knowledge, the existence of this polycarbonyl species has been explicitly rejected by literature. In Si(CO)<sub>2</sub>, silicon dicarbonyl, the CO are dative-bonded to silicon; Si(CO)<sub>4</sub>, silicon tetracarbonyl, may be viewed as a resonance between the extreme configurations (CO)<sub>2</sub>Si + 2CO and 2CO + Si(CO)<sub>2</sub>; while Si[C<sub>2</sub>O<sub>2</sub>], *c*-silicodiketone, is somewhat similar to the compounds formed by silicon and ethylene. A detailed orbital analysis has revealed how in the Si(CO)<sub>2</sub> complex the nature of the Si bonding is consistent with the use of  $sp^2$ -hybridized orbitals on silicon. For the ( $S = 0$ ) silicon dicarbonyl complex the  $\sigma$  interaction between Si and (CO)<sub>2</sub> framework due to electron

donation from carbonyl to silicon accounts for most of the bond strength, but  $\pi$  backdonation from silicon  $3p_\pi$  to empty  $2\pi^*$  CO orbitals takes place. For the ( $S = 1$ ) *c*-silicodiketone complex additional  $\sigma$  interaction occurs between the two C atoms and  $\pi$  backdonation from silicon to  $2\pi^*$  CO orbitals occurs.

The Si bonding in Si(CO)<sub>4</sub> is consistent with the use of  $sp^2d$ -hybridized orbitals on silicon. This complex represents a clear example of  $d$  orbital involvement in the formation of a four-coordinated planar silicon.

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